

Zeeman effect implementation in ARTS

Nikolay Koulev

Institute for Environmental Physics

nkoulev@uni-bremen.de

<http://www.sat.uni-bremen.de>

Bredbeck

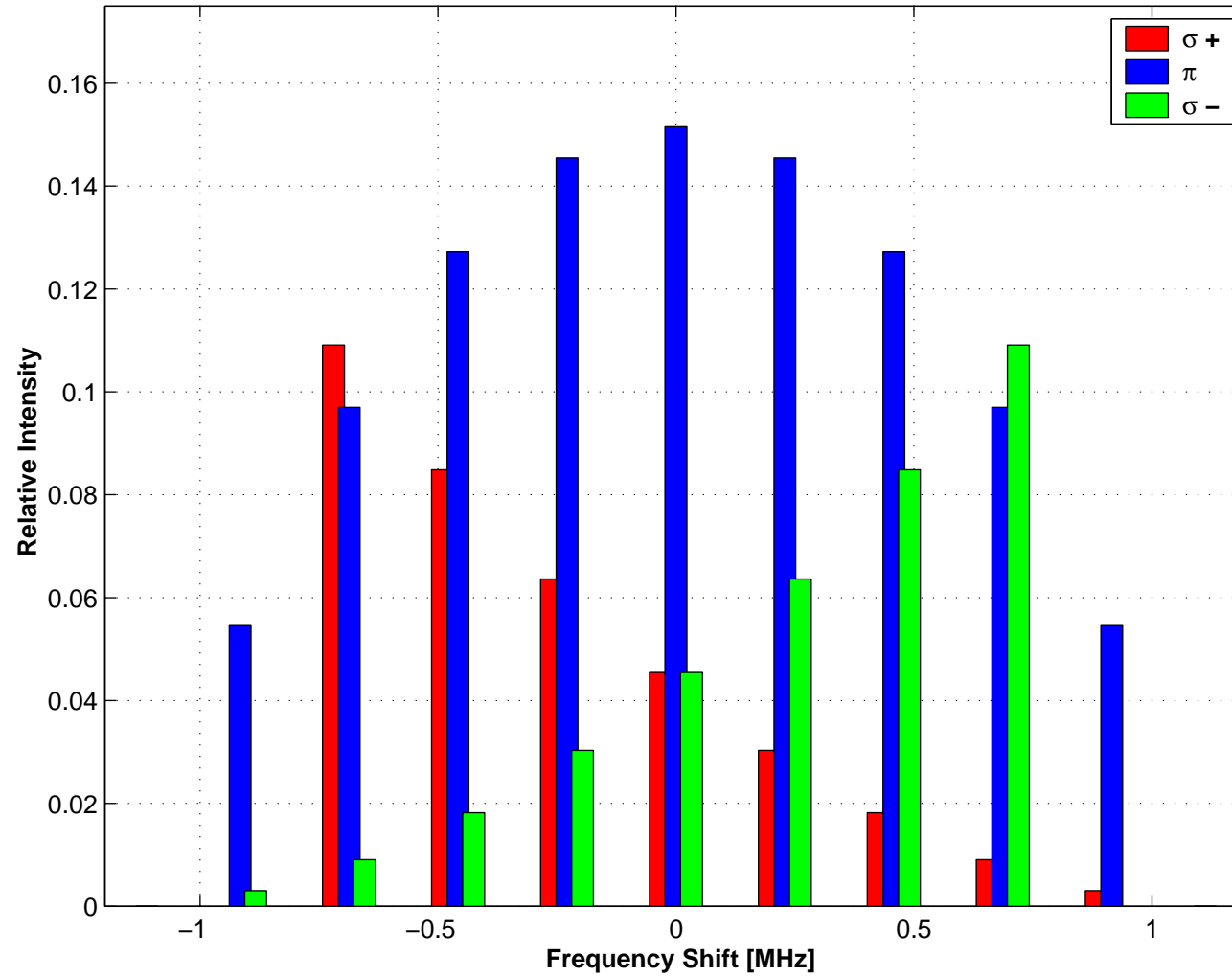
8. July 2003



-
- ▶ Introduction
 - ▶ ARTS RTE modelling and the Zeeman effect
 - ▶ Radiation matrices and polarization bases
 - ▶ Implementation issues

Zeeman effect

5- O2 line @ 60.306061 GHz at 95 km by 35000 nT



Non-scattering Vector RTE in ARTS

$$\frac{d}{ds}\vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu) = -\hat{\mathbf{K}}(\vec{\mathbf{n}}, \nu) \cdot \vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu) + \vec{\mathbf{a}}(\vec{\mathbf{n}}, \nu)B(\nu)$$

- $\vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu)$ - 4D Stokes vector
- $\hat{\mathbf{K}}(\vec{\mathbf{n}}, \nu)$ - 4x4 Extinction coefficient matrix
- $\vec{\mathbf{a}}(\vec{\mathbf{n}}, \nu)$ - 4D Absorption coefficient vector
- $\vec{\mathbf{n}}$ - propagation vector
- $B(\nu)$ - Planck function

Zeeman effect RTE

$$\frac{d}{ds}\hat{\mathbf{R}}(\vec{\mathbf{n}}, \nu) = -[\hat{\mathbf{G}}(\nu) \cdot \hat{\mathbf{R}}(\vec{\mathbf{n}}, \nu) + \hat{\mathbf{R}}(\vec{\mathbf{n}}, \nu) \cdot \hat{\mathbf{G}}^*(\nu)] \\ + B(\nu)[\hat{\mathbf{G}}(\nu) + \hat{\mathbf{G}}^*(\nu)]$$

- $\hat{\mathbf{R}}(\vec{\mathbf{n}}, \nu)$ - 2x2 Radiation Intensity (Brightness temperature*) Matrix
- $\hat{\mathbf{G}}(\nu)$ - 2x2 complex Propagation tensor

$$\hat{\mathbf{G}}(\nu) = i \frac{2\pi\nu}{c} [\hat{\mathbf{I}} + \hat{\mathbf{M}}_{AN}(\nu)]$$

*then using the medium physical temperature T instead of $B(\nu)$

ARTS Zeeman effect approach

$$\frac{d}{ds}\vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu) = -\hat{\mathbf{K}}_Z(\vec{\mathbf{n}}, \nu) \cdot \vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu) + \vec{\mathbf{a}}_Z(\vec{\mathbf{n}}, \nu)B(\nu)$$

where

$$\hat{\mathbf{K}}_Z(\vec{\mathbf{n}}, \nu) = \hat{f}(\hat{\mathbf{G}}(\nu))$$

$$\vec{\mathbf{a}}_Z(\vec{\mathbf{n}}, \nu) = \vec{g}(\hat{\mathbf{G}}(\nu))$$

Radiation matrices

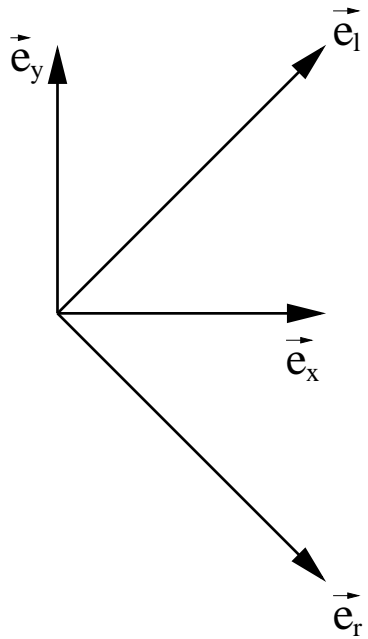
$$\hat{\mathbf{R}}(\vec{\mathbf{n}}, \nu) \sim \hat{\mathbf{G}}(\nu) \cdot \left\langle \vec{\mathbf{H}}(\vec{\mathbf{n}}, \nu) \otimes \vec{\mathbf{H}}^*(\vec{\mathbf{n}}, \nu) \right\rangle$$

- $\vec{\mathbf{H}}(\vec{\mathbf{n}}, \nu)$ - 2D complex field vector of the propagated wave.

$$\hat{\mathbf{C}}(\vec{\mathbf{n}}, \nu) = \left\langle \vec{\mathbf{H}}(\vec{\mathbf{n}}, \nu) \otimes \vec{\mathbf{H}}^*(\vec{\mathbf{n}}, \nu) \right\rangle$$

- $\hat{\mathbf{C}}(\vec{\mathbf{n}}, \nu)$ - complex Coherency matrix of the propagated wave.

Polarization bases



- \vec{e}_x, \vec{e}_y - linear polarization basis.
- \vec{e}_l, \vec{e}_r - circular polarization basis.

$$\begin{pmatrix} \vec{e}_l \\ \vec{e}_r \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \\ 1 & -i \end{pmatrix} \cdot \begin{pmatrix} \vec{e}_x \\ \vec{e}_y \end{pmatrix}$$

Coherency Matrix and Stokes Vector

In an arbitrary (\vec{e}_a, \vec{e}_b) polarization basis:

$$\hat{\mathbf{C}}^{(ab)} = \frac{1}{2} \sum_{i=0}^3 \mathbf{I}_i^{(ab)} \hat{\sigma}_i$$

where

$$\hat{\sigma}_0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \hat{\sigma}_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \hat{\sigma}_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \hat{\sigma}_3 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

and $\mathbf{I}_i^{(ab)}$ are the Stokes vector components.

Solution for $\hat{\mathbf{K}}_Z(\vec{\mathbf{n}}, \nu)$ and $\vec{\mathbf{a}}_Z(\vec{\mathbf{n}}, \nu)$

In a circular $(\vec{\mathbf{e}}_l, \vec{\mathbf{e}}_r)$ polarization basis:

$$\hat{\mathbf{K}}_Z^C = \frac{1}{2} \begin{pmatrix} A & B & C & 0 \\ B & A & 0 & -D \\ C & 0 & A & E \\ 0 & D & -E & A \end{pmatrix} \quad \vec{\mathbf{a}}_Z^C = \begin{pmatrix} A \\ B \\ C \\ 0 \end{pmatrix}$$

where $A = 2\text{Re}[G_{11} + G_{22}]$, $B = -2\text{Re}[G_{22} - G_{11}]$, $C = 2\text{Re}[G_{12} + G_{21}]$,
 $D = 2\text{Im}[G_{12} + G_{21}]$ and $E = -2\text{Im}[G_{22} - G_{11}]$.

Stokes vector transformation

$$\vec{\mathbf{I}}_C(\vec{\mathbf{n}}, \nu) = \hat{\mathbf{V}}_C \cdot \vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu)$$

where

$$\hat{\mathbf{V}}_C = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \end{pmatrix}$$

and $\vec{\mathbf{I}}_C(\vec{\mathbf{n}}, \nu)$, $\vec{\mathbf{I}}(\vec{\mathbf{n}}, \nu)$ - Stokes vector in circular and linear polarization basis.

ARTS final Zeeman effect RTE

In a linear (\vec{e}_x, \vec{e}_y) polarization basis:

$$\frac{d}{ds} \vec{I}(\vec{n}, \nu) = - \overbrace{\hat{V}_C^{-1} \cdot \hat{K}_Z^C(\vec{n}, \nu) \cdot \hat{V}_C}^{\hat{K}_Z(\vec{n}, \nu)} \cdot \vec{I}(\vec{n}, \nu) + \underbrace{\hat{V}_C^{-1} \cdot \vec{a}_Z^C(\vec{n}, \nu)}_{\vec{a}_Z(\vec{n}, \nu)} B(\nu)$$

Additional implementation issues

The fully fledged Zeeman part in ARTS will utilize:

- all the pressure and temperature profiles provided by `arts-data` and `arts-xml-data` packages
- the IGRF2000 profiles of the global magnetic field for Epoch 2000-2005 with possibility of an update
- (possibly) the line-mixing mechanism between the split lines

References

Lenoir, W. B., *Propagation of partially polarized waves in a slightly anisotropic medium*, **Journal of Applied Physics**, Vol. 38, p. 5283-5290, 1967

J. Pardo, M. Ridal, D. Murtagh and J. Chernicharo, *Microwave teperature and pressure measurements with the Odin satellite: I. Observational method*, **Canadian Journal of Physics**, Volume 80, Number 4, April 2002
ISSN 1208-6045

Christian Brossau, *Fundamentals of polarized light: a statistical approach*, 1998, John Wiley & Sons, Inc., ISBN 0-471-14302-2

William Swindell(Ed.), *Benchmark Papers in Optics; v.1, Polarized light*, 1975, Dowden, Hutchinson & Ross, Inc., ISBN 0-470-83997-X